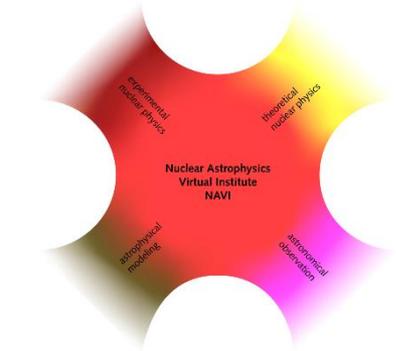


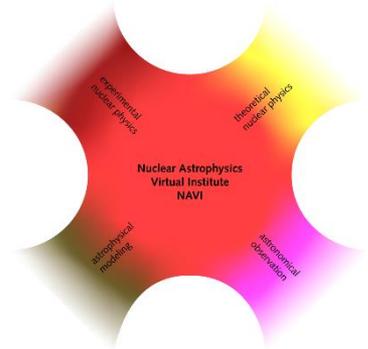
# Improving Microphysics of Neutrino-Nucleon Interactions in Supernovae

MICRA 2015 – August 21– Stockholm  
Andreas Lohs (Univ. Basel)



# (Almost) Improving Microphysics of Neutrino-Nucleon Interactions in Supernovae

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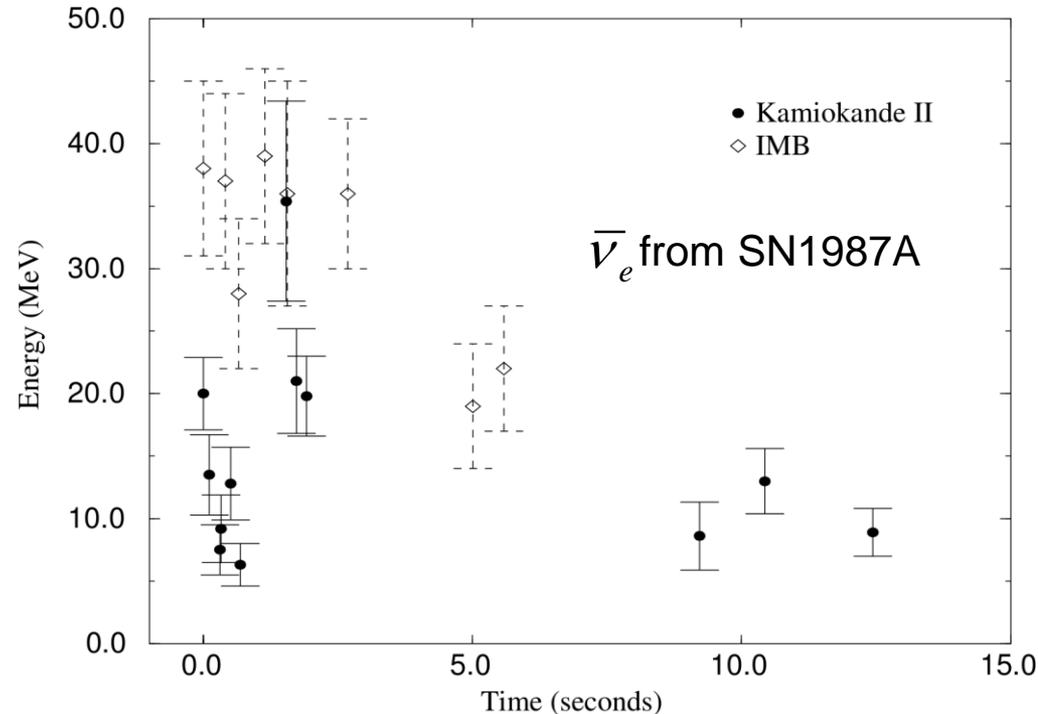
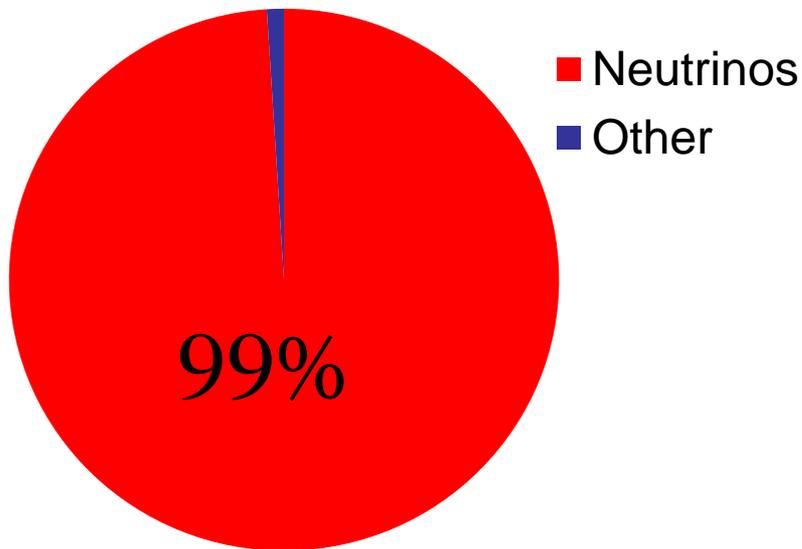
# Improving Microphysics of Neutrino-Nucleon Interactions in Supernovae (in Medium)

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# Neutrinos in Supernovae

Core collapse supernovae release huge amount of energy

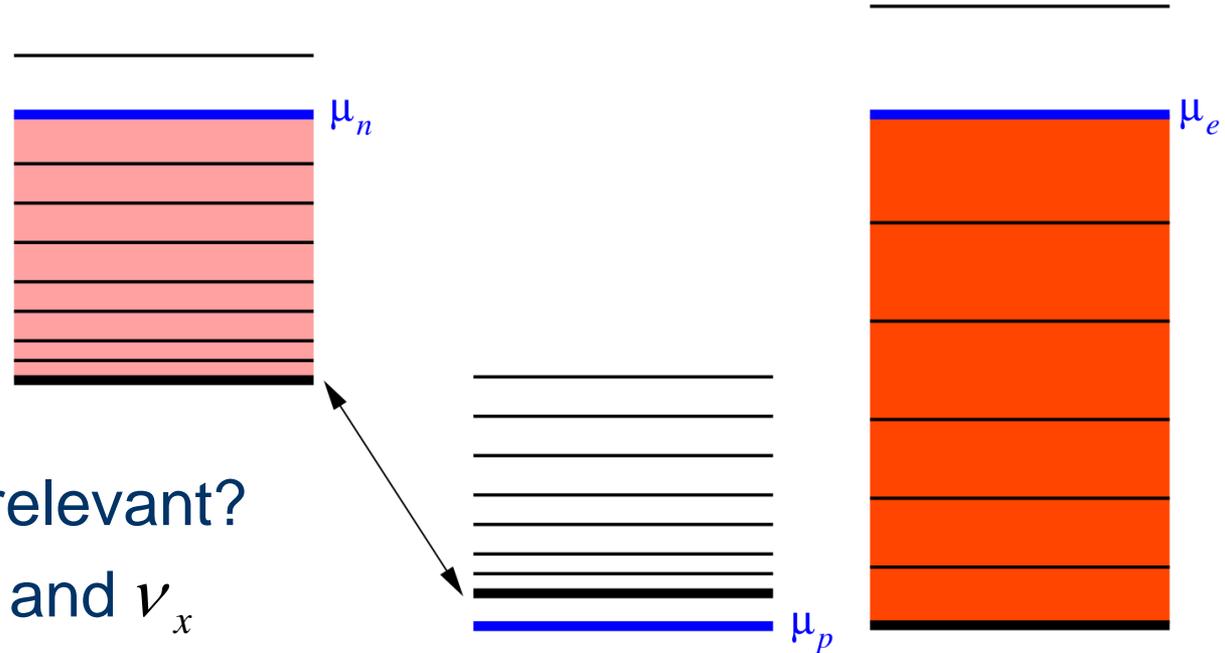
## Supernova energy



As neutrinos affect many aspects of Snc, what are main questions of neutrino transport?

# Uncertainties in Neutrino Physics

What is the correct Equation of state?



Which reactions are relevant?

- Not obvious for  $\bar{\nu}_e$  and  $\nu_x$
- Answer may vary for different SNe

**How to compute neutrino interactions?**

- inelasticity, relativity, medium effects, weak magnetism ...

# Mean Free Path for Neutrino Absorption

## Elastic Approximation

- Lowest order expression for nonrelativistic nucleons
- Analytic formula for  $\lambda(E_\nu)$
- Can be corrected to include recoil, weak magnetism, ...

## Nucleons as quasi-free fermions

- Relativistic kinematics, „full“ matrix element, no correlations
- Mostly 2-D numerical integrals to obtain  $\lambda(E_\nu)$

## Structure function from RPA / Linear response theory

- Fully consistent with RMF-EOS, correlations (can be) included
- Requires 3-D numerical integrals to obtain  $\lambda(E_\nu)$

# Elastic Approximation for Neutrino Absorption

Mean-free path for (quasi-) free particles:

$$\lambda(E_\nu)^{-1} \sim \int d^3 p_e [1 - f_e(E_e)] \int d^3 p_n \int d^3 p_p \frac{\langle |M|^2 \rangle}{16 E_\nu E_n E_e E_p} f_n(E_n) [1 - f_p(E_p)] \delta^4$$

Assume non-relativistic nucleons and elastic collision:

$$E_{n,p} \simeq m_{n,p} \Rightarrow \frac{\langle |M|^2 \rangle}{16 E_\nu E_n E_e E_p} \simeq G_A^2 (3 - x) + G_V^2 (1 + x)$$

$$E_n - E_p \simeq m_n - m_p + U_n - U_p$$

Mean-free path reduces to

$$\lambda(E_\nu)^{-1} \sim (3G_A^2 + G_V^2) (E_\nu + \Delta m + \Delta U)^2 [1 - f_e(E_\nu + \Delta m + \Delta U)] \frac{n_n - n_p}{1 - \exp[(\eta_p - \eta_n)/T]}$$

# Recoil and Weak Magnetism Corrections

[Horowitz, PRD 65 (2002) 043001] pointed out:

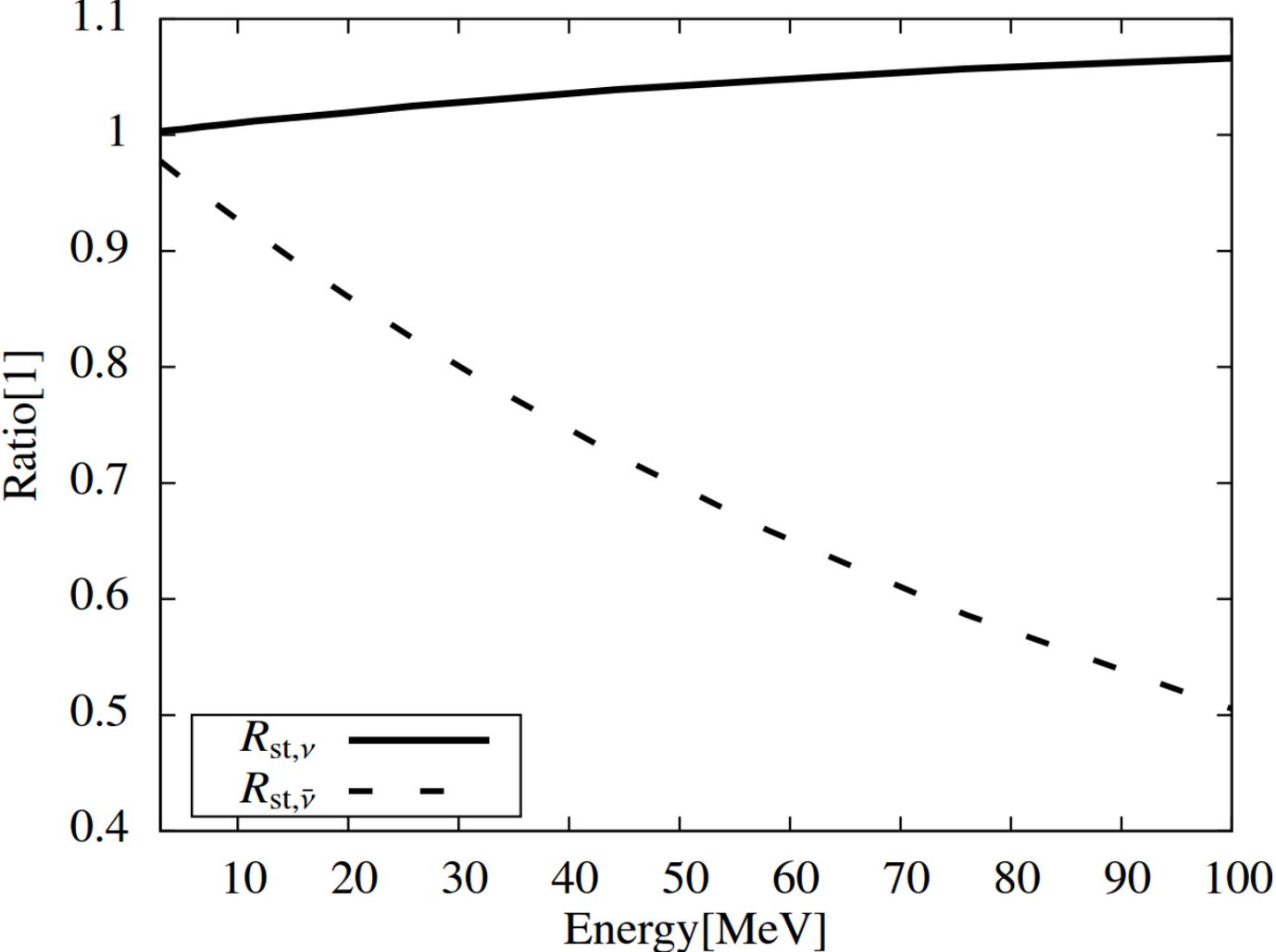
- “Elastic Approximation“ is more simplified than necessary
- Kinematics/Recoil can be treated relativistically

$$E_n = m_n \Rightarrow E_e = \frac{E_\nu}{1 + \frac{E_\nu}{m_n} (1 - x)}$$

- Include in phase space factor and matrix element
- Gives rise to analytic correction factor for cross-section

$$R = \left\{ G_V^2 \left( 1 + 4e + \frac{16}{3}e^2 \right) + 3G_A^2 \left( 1 + \frac{4}{3}e \right)^2 \pm 4G_A (G_V + F_2) e \left( 1 + \frac{4}{3}e \right) + \frac{8}{3}G_V F_2 e^2 + \frac{1}{3}F_2^2 e^2 (5 + 2e) \right\} / \left[ (1 + 2e)^3 (G_V^2 + 3G_A^2) \right]$$

# Correction Factor for Cross-Section



# Improvement: Consider Mass and Potential Differences

- Masses and strong interaction potentials of nucleons differ
- At large densities effective masses decrease

$$E_e = \frac{E_\nu + \frac{M_*^2 - m_p^{*2}}{2M_*}}{1 + \frac{E_\nu}{M_*} (1 - x)}$$

$$M_* = m_n^* + U_n - U_p$$

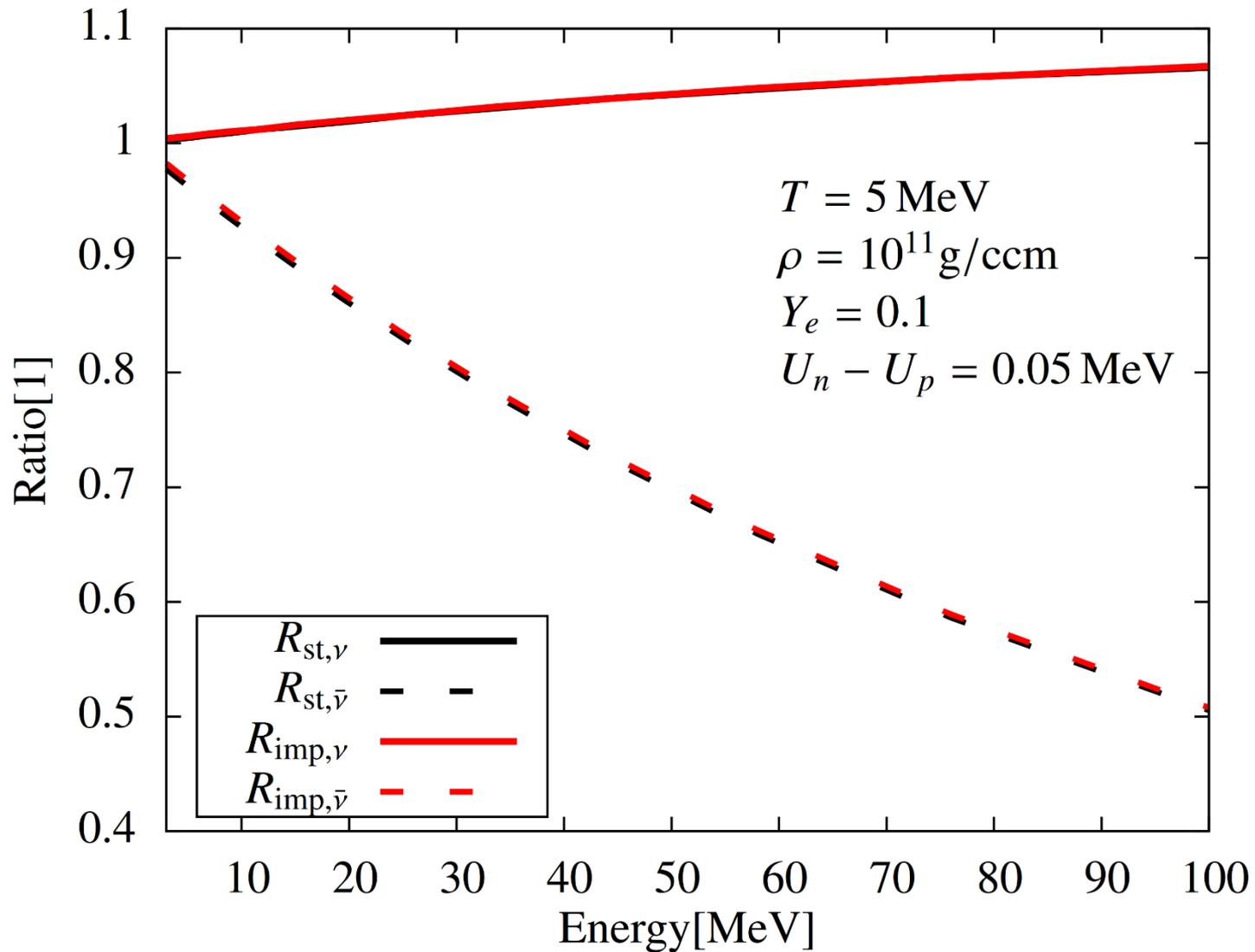
- Analytic correction factor can still be derived the same way
- In the matrix element, additional terms can be included
- For neutrino scattering, only difference is exchange of rest mass with effective mass

# Improved Correction Factor

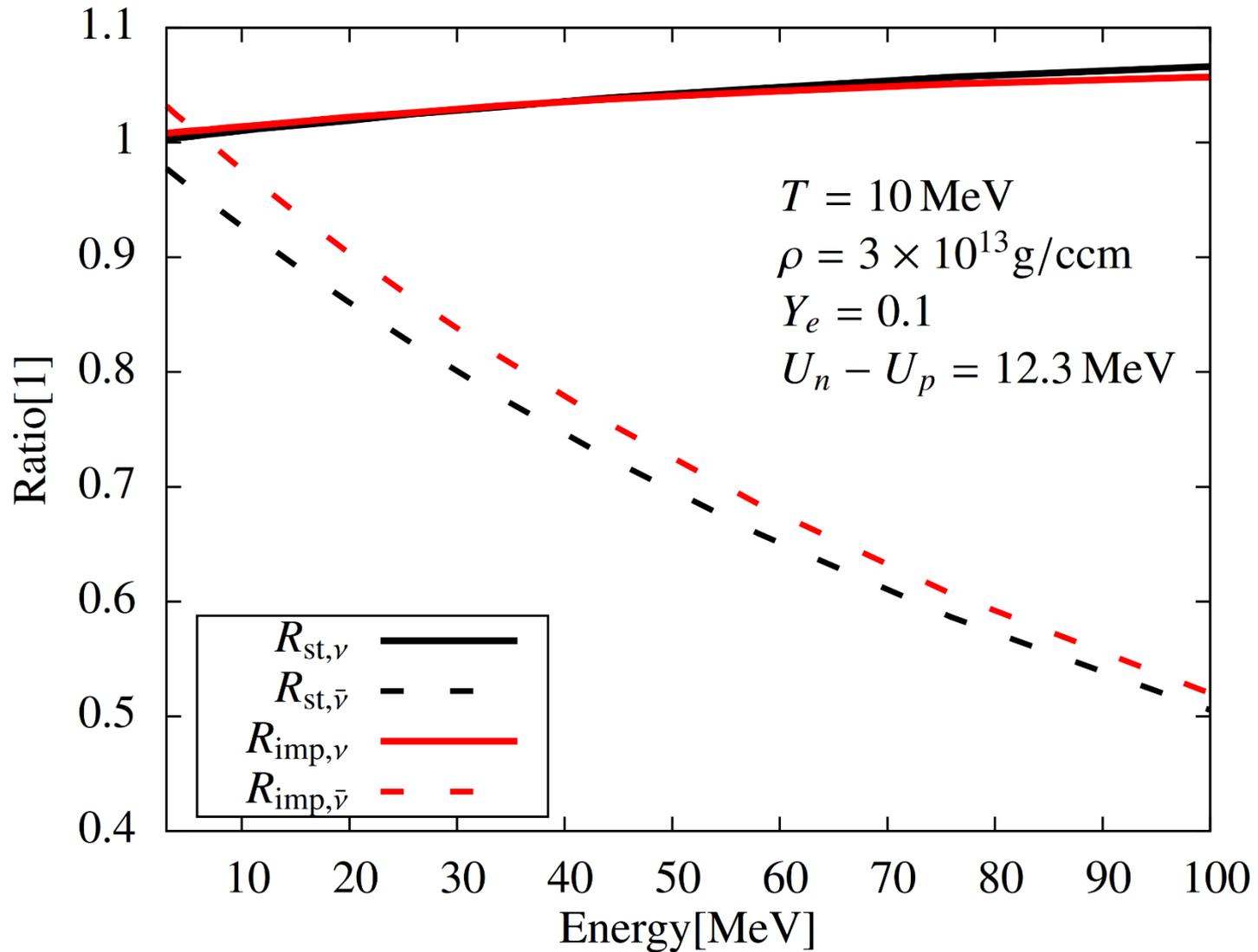
$$\begin{aligned}
 R = & \left\{ G_V^2 \left[ 1 + 4e_* + \frac{16}{3}e_*^2 + \frac{4}{3}e_*\xi + \left( 1 + \frac{2}{3}e_* \right) (\xi - q_*) \right] \right. \\
 & + G_A^2 \left[ 3 + 8e_* + \frac{16}{3}e_*^2 - \frac{4}{3}e_*\xi - \left( 1 + \frac{2}{3}e_* \right) (\xi + q_*) \right] \\
 & \pm G_A \left[ G_V + F_2 \frac{M_*}{m_N} \left( 1 - \frac{\xi}{2} \right) \right] \left[ 4e_* + \frac{16}{3}e_*^2 + q_* \left( 2 + \frac{4}{3}e_* \right) \right] \\
 & + G_V F_2 \frac{M_*}{m_N} \left[ \left( 1 + \frac{q_*}{e_*} - \frac{\xi}{2} \right) \frac{8}{3}e_*^2 + \xi q_* \left( 1 + 2e_* + \frac{4}{3}e_*^2 \right) \right] \\
 & \left. + F_2^2 \frac{M_*^2}{m_N^2} \left[ \frac{5}{3}e_*^2 + \frac{2}{3}e_*^3 + \left( \frac{1}{2} + e_* \right) \tilde{A} + \left( \frac{1}{2} + \frac{1}{3}e_* \right) \tilde{B} + \frac{2}{3}e_* \tilde{C} \right] \right\} \\
 & / \left[ (1 + 2e)^3 (G_V^2 + 3G_A^2) \right]
 \end{aligned}$$

$$\xi = \frac{\Delta m^* + \Delta U}{M_*}, \quad q = \frac{m_n^{*2} - m_p^{*2}}{2M_*^2}, \quad q_* = \frac{M_*^2 - m_p^{*2}}{2M_*^2}$$

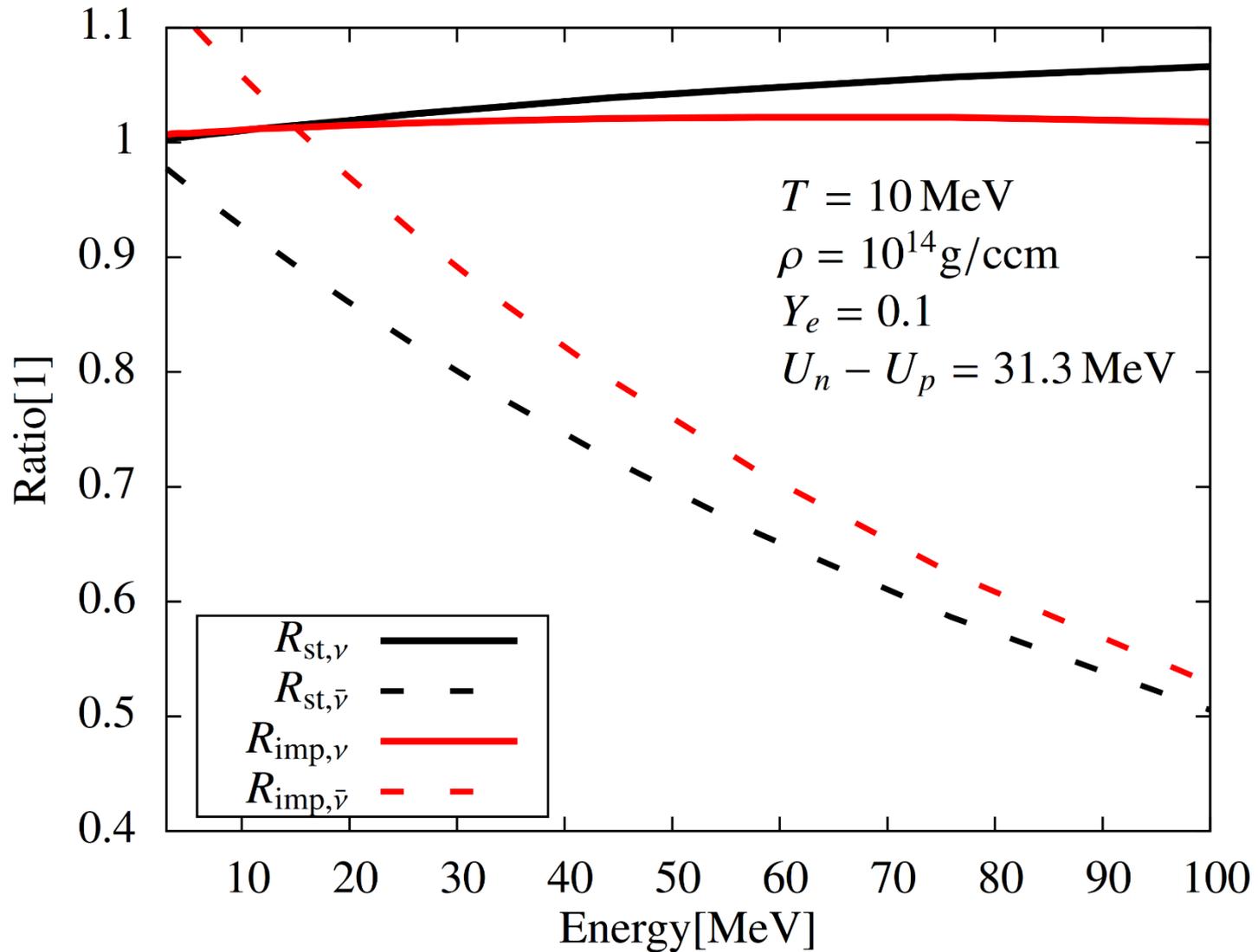
# Improved Correction Factor at Low Densities



# Improved Correction Factor at High Densities

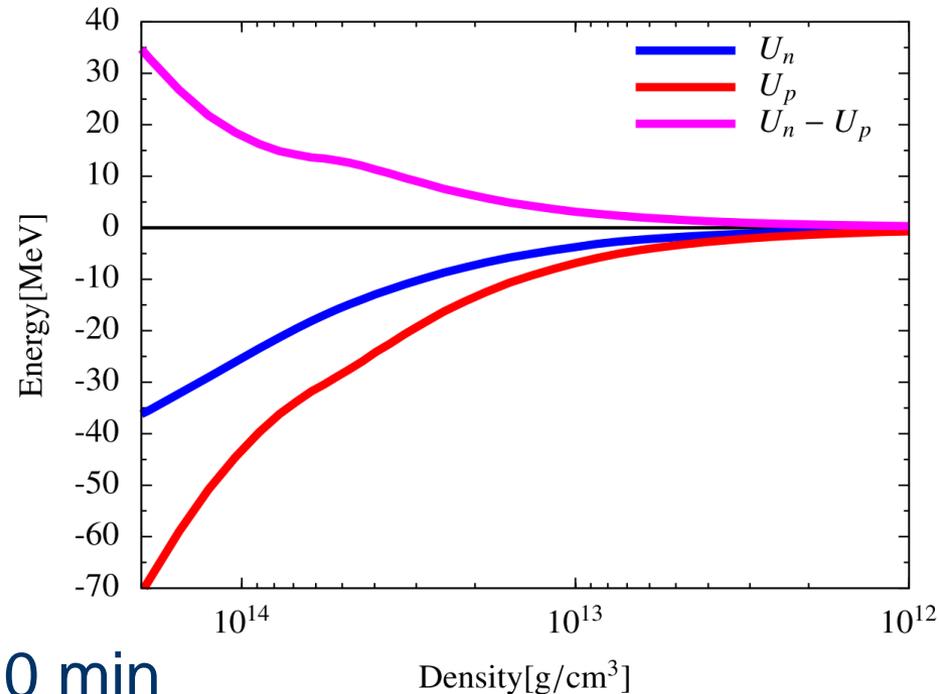
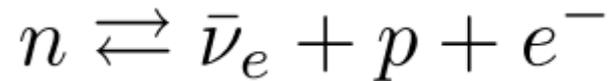
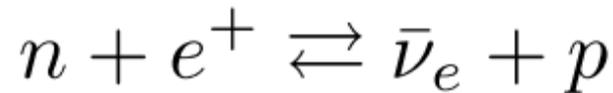


# Improved Correction Factor at High Densities



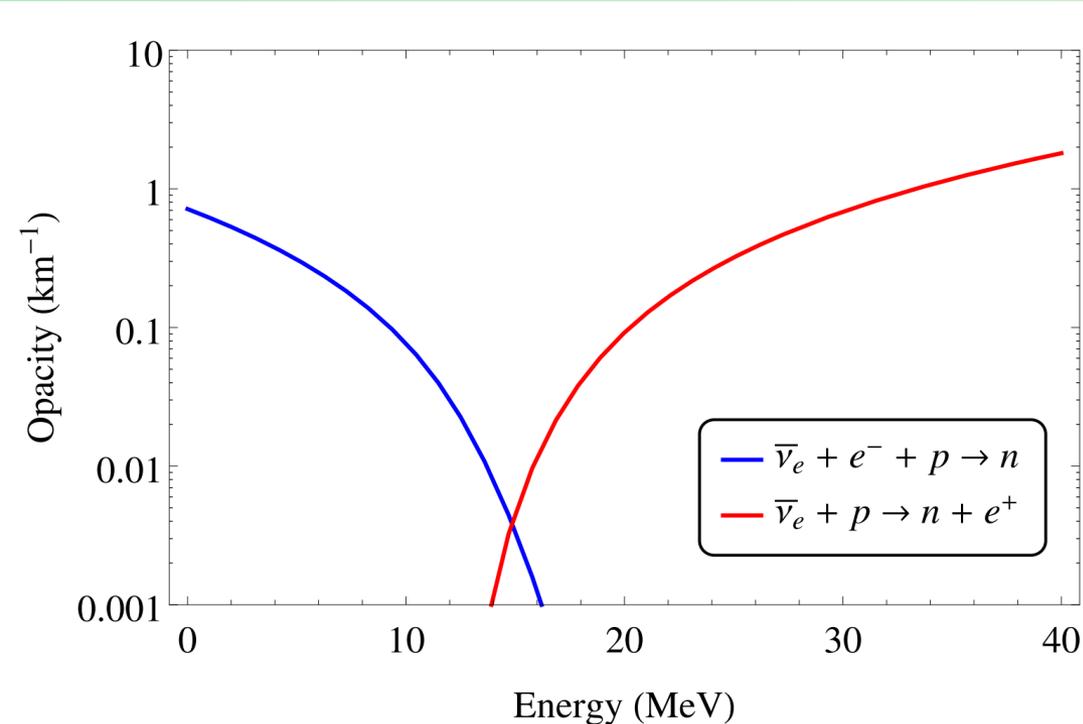
# Neutron decay at high density

- Low energy  $\bar{\nu}_e$  cannot be absorbed on protons or produced from positron capture for large  $U_n - U_p$

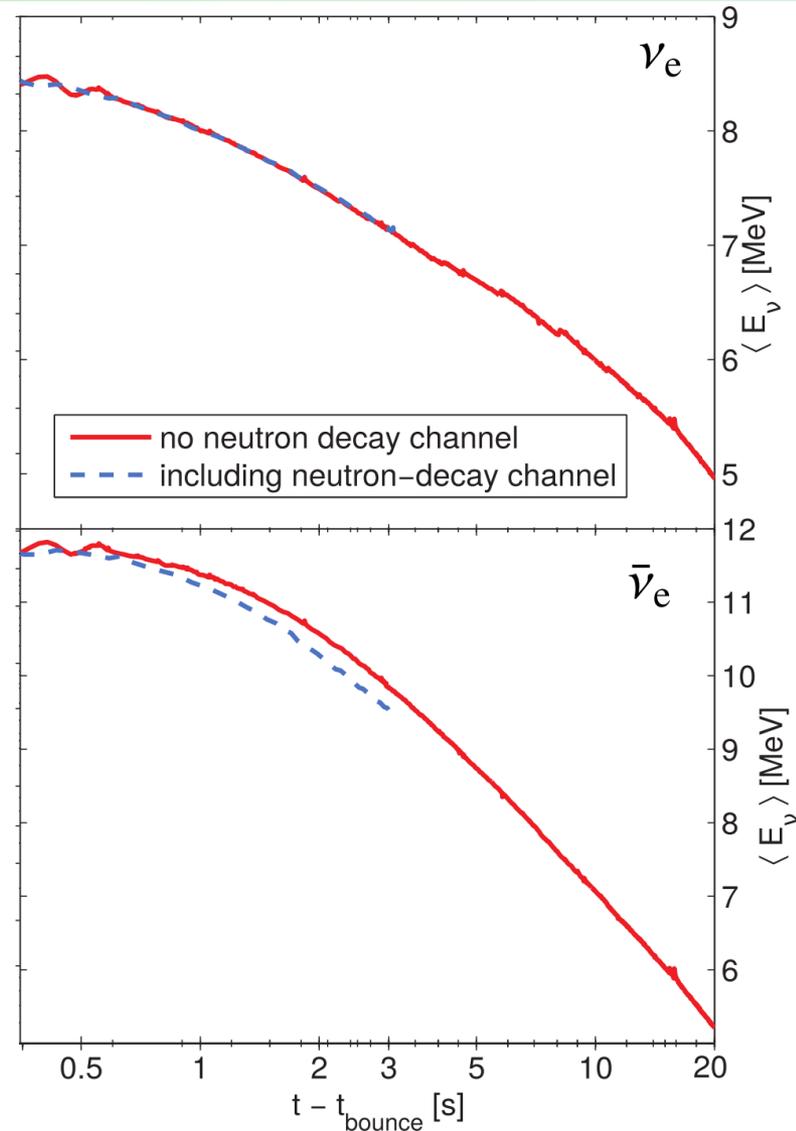


- Neutron lifetime in vacuum  $\sim 10$  min
- Strong interaction increases Q-value for nucleon conversion  
→ Decay rate raises

# Neutron decay at high density



- Decreasing average energy of  $\bar{\nu}_e$
  - Spectral change will affect nucleosynthesis yields
- Rising  $Y_e$



# Elastic Approximation and Corrections for Neutron Decay

Elastic approximation for neutron decay similar to absorption

$$\lambda(E_\nu)^{-1} \sim (3G_A^2 + G_V^2) (\Delta m + \Delta U - E_\nu)^2 f_e(\Delta m + \Delta U - E_\nu) \frac{n_p - n_n}{1 - \exp[(\eta_n - \eta_p)/T]}$$

Kinematic relation for inverse neutron decay, assuming proton at rest

$$E_e = \frac{-E_{\bar{\nu}} + \frac{M_*^2 - M_f^2}{2M_*}}{1 + \frac{E_{\bar{\nu}}}{M_*} (1 - x)}$$

Physical meaning only for different nucleon masses and/or potentials

# Summary and Conclusion

## Correction Factors for neutrino-nucleon interactions

- Can be extended to include strong interaction potentials and effective masses
- Shifts Correction factors at high densities
- Correction factor for neutron decay requires improvement to have physical meaning

## Outlook:

- Correction factor for neutron decay
- Include finite lepton masses
- Include momentum dependence of coupling constants